

# Lec 6

Story of

$$\dot{x}(t) = A(t)x(t) + f(t).$$

6.1

We are looking at eq<sup>s</sup> of the form

$$\dot{x}(t) = A(t)x(t) + f(t). \quad (*)$$

Problem is to solve  $(*)$

Special case 1 ( $A(t) \equiv 0$ )

$$\dot{x}(t) = f(t)$$

$$x(t) = ~~x(t_0)~~ x(t_0) + \int_{t_0}^t f(\sigma) d\sigma$$

↑  
initial condition  
at  $t = t_0$ .

Special case 2 ( $f(t) = 0$ )

$$\dot{x}(t) = A(t)x(t).$$

$$x(t) = \Phi(t, t_0)x(t_0).$$

↑  
transition matrix.

Define a new variable

$$z(t) = \phi(t_0, t) x(t)$$

Remark:  $\phi(t, t_0) = \phi(t_0, t)^{-1}$

Lemma:  $\dot{\phi}(t_0, t) = -\phi(t_0, t) A(t)$ .

$$\phi(t, t_0) \phi(t_0, t) = I$$

Taking the derivative w.r.t.  $t$  we have

$$\frac{d}{dt} \phi(t, t_0) \phi(t_0, t) + \phi(t, t_0) \frac{d}{dt} \phi(t_0, t) = 0$$

$$\Rightarrow A(t) \underbrace{\phi(t, t_0) \phi(t_0, t)}_{= I} + \phi(t, t_0) \frac{d}{dt} \phi(t_0, t) = 0$$

$$\Rightarrow \frac{d}{dt} \phi(t_0, t) = -\phi(t, t_0)^{-1} A(t)$$

$$= -\phi(t_0, t) A(t)$$

6.3

It follows that

$$\dot{z}(t) = \dot{\phi}(t_0, t) x(t) + \phi(t_0, t) \dot{x}(t)$$

$$= -\phi(t_0, t) A(t) x(t) +$$

$$\phi(t_0, t) A(t) x(t) + \phi(t_0, t) f(t)$$

$$= \phi(t_0, t) f(t).$$

Thus

$$\dot{z}(t) = \phi(t_0, t) f(t)$$

This implies

$$z(t) = z(t_0) + \int_{t_0}^t \phi(t_0, \sigma) f(\sigma) d\sigma.$$

$$\Rightarrow \phi(t_0, t) x(t) = x(t_0) + \int_{t_0}^t \phi(t_0, \sigma) f(\sigma) d\sigma$$

6.4

Thus we have

$$x(t) = \Phi(t_0, t)^{-1} x(t_0) +$$

$$\int_{t_0}^t \Phi(t_0, t)^{-1} \Phi(t_0, \sigma) f(\sigma) d\sigma$$

$$= \Phi(t, t_0) x(t_0) +$$

$$\int_{t_0}^t \Phi(t, t_0) \Phi(t_0, \sigma) f(\sigma) d\sigma.$$

Theorem 6.1

(Variation of constants formula)

Let  $\phi(t, t_0)$  be the state transition matrix for  $\dot{x}(t) = A(t)x(t)$ , then the unique sol<sup>n</sup> of

$$\dot{x}(t) = A(t)x(t) + f(t)$$

$$x(t_0) = x_0$$

is given by

$$x(t) = \phi(t, t_0)x_0 + \int_{t_0}^t \phi(t, \sigma)f(\sigma)d\sigma.$$

Corollary 6.2

The solution of the inhomogeneous linear constant equation  $\dot{x} = Ax + f(t)$ ;

$x(0) = x_0$  is given by

$$x(t) = e^{At}x_0 + \int_{t_0}^t e^{A(t-\sigma)}f(\sigma)d\sigma.$$

Example(Newton's 2<sup>nd</sup> law)

For a unit mass

$$\ddot{x} = f(t).$$

$\uparrow$                      $\uparrow$   
 acc<sup>n</sup>                force

Define  $x_1 = x$ ,  $x_2 = \dot{x}$ , we have

$$\begin{pmatrix} \dot{x}_1 \\ \dot{x}_2 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} + \begin{pmatrix} 0 \\ f(t) \end{pmatrix}$$

$$\begin{pmatrix} x_1(t) \\ x_2(t) \end{pmatrix} = \begin{pmatrix} 1 & t \\ 0 & 1 \end{pmatrix} \begin{pmatrix} x_1(0) \\ x_2(0) \end{pmatrix} + \int_0^t \begin{pmatrix} 1 & t-\sigma \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 0 \\ f(\sigma) \end{pmatrix} d\sigma$$

Hence

$$x_1(t) = x_1(0) + t x_2(0) + \int_0^t (t-\sigma) f(\sigma) d\sigma$$

$$x_2(t) = x_2(0) + \int_0^t f(\sigma) d\sigma$$

In the original variables we write

(6.7)

$$x(t) = x(0) + t \dot{x}(0) + \int_0^t (t-\sigma) f(\sigma) d\sigma .$$

↑  
Integral form of Newton's Law



Example (The satellite problem)

$$\dot{\underline{x}} = A \underline{x} + B u$$

$$A = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 3\omega^2 & 0 & 0 & 2\omega \\ 0 & 0 & 0 & 1 \\ 0 & -2\omega & 0 & 0 \end{pmatrix} ; B = \begin{pmatrix} 0 & 0 \\ 1 & 0 \\ 0 & 0 \\ 0 & 1 \end{pmatrix}$$

The state transition matrix  $\phi(t, t_0)$  is given by

$$\left( \begin{array}{cccc} 4 - 3\cos\omega(t-t_0) & \frac{\sin\omega(t-t_0)}{\omega} & 0 & \frac{2(1-\cos\omega(t-t_0))}{\omega} \\ 3\omega\sin\omega(t-t_0) & \cos\omega(t-t_0) & 0 & 2\sin\omega(t-t_0) \\ 6(-\omega(t-t_0) + \sin\omega(t-t_0)) & \frac{-2(1-\cos\omega(t-t_0))}{\omega} & 1 & \frac{-3\omega(t-t_0) + 4\sin\omega(t-t_0)}{\omega} \\ 6\omega(-1 + \cos\omega(t-t_0)) & -2\sin\omega(t-t_0) & 0 & -3 + 4\cos\omega(t-t_0) \end{array} \right)$$

6.9

Starting from  $x(0) = x_0$ , we have

$$x(t) = \phi(t, 0) x_0 + \int_0^t \phi(t, \sigma) B(\sigma) d\sigma$$

where

$$\phi(t, \sigma) B(\sigma) = .$$

$$\begin{bmatrix} \frac{\sin \omega(t-\sigma)}{\omega} & \frac{2(1-\cos \omega(t-\sigma))}{\omega} \\ \cos \omega(t-\sigma) & 2 \sin \omega(t-\sigma) \\ \frac{-2(1-\cos \omega(t-\sigma))}{\omega} & \frac{-3\omega(t-\sigma) + 4 \sin \omega(t-\sigma)}{\omega} \\ -2 \sin \omega(t-\sigma) & -3 + 4 \cos \omega(t-\sigma) \end{bmatrix}$$

An example that we cannot live ①  
without

The Euler Equations for the angular velocities of a rigid body are

$$I_1 \dot{\omega}_1 = (I_2 - I_3) \omega_2 \omega_3 + u_1$$

$$I_2 \dot{\omega}_2 = (I_3 - I_1) \omega_1 \omega_3 + u_2$$

$$I_3 \dot{\omega}_3 = (I_1 - I_2) \omega_1 \omega_2 + u_3$$

$\omega = \begin{pmatrix} \omega_1 \\ \omega_2 \\ \omega_3 \end{pmatrix}$  is the angular velocity in a body fixed co-ordinate system coinciding with the principal axes

$u = \begin{pmatrix} u_1 \\ u_2 \\ u_3 \end{pmatrix}$  is the applied torque.

$I_1, I_2, I_3$  are the principal moment of inertia.

(2)

If  $I_1 = I_2$  we call the body symmetrical. In this case, for  $u=0$  we have

$$\dot{\omega}_3 = 0 \Rightarrow \omega_3 = \omega_0 \text{ (a constant).}$$

Hence

$$\dot{\omega}_1 = \left[ \left( \frac{I_2 - I_3}{I_2} \right) \omega_0 \right] \omega_2$$

$$\dot{\omega}_2 = \left[ \left( \frac{I_3 - I_2}{I_2} \right) \omega_0 \right] \omega_1$$

Define  $\alpha = \left( \frac{I_2 - I_3}{I_2} \right) \omega_0$ , it follows that

$$\boxed{\begin{array}{l} \dot{\omega}_1 = \alpha \omega_2 \\ \dot{\omega}_2 = -\alpha \omega_1 \end{array}} \quad \begin{bmatrix} \omega_1(t) \\ \omega_2(t) \end{bmatrix} = \begin{pmatrix} \cos \alpha t & \sin \alpha t \\ -\sin \alpha t & \cos \alpha t \end{pmatrix} \begin{pmatrix} \omega_1(0) \\ \omega_2(0) \end{pmatrix}$$

Choose  $\omega_1(0)=1, \omega_2(0)=0$ , we have

(3)

$$\omega_3 = \omega_0, \omega_2 = -\sin \alpha t, \omega_1 = \cos \alpha t$$

as a sol<sup>n</sup> to the symmetrical Euler eq<sup>s</sup>

Define

$$\bar{\omega}_1 = \omega_1 - \cos \alpha t$$

$$\bar{\omega}_2 = \omega_2 + \sin \alpha t$$

$$\bar{\omega}_3 = \omega_3 - \omega_0$$

We have

$$\dot{\bar{\omega}}_3 = \dot{\omega}_3 = 0$$

$$\begin{aligned} \dot{\bar{\omega}}_1 &= \dot{\omega}_1 + \alpha \sin \alpha t = \alpha \omega_2 + \alpha \sin \alpha t \\ &= \alpha \bar{\omega}_2 \end{aligned}$$

$$\begin{aligned} \dot{\bar{\omega}}_2 &= \dot{\omega}_2 + \alpha \cos \alpha t = -\alpha (\omega_1 - \cos \alpha t) \\ &= -\alpha \bar{\omega}_1 \end{aligned}$$

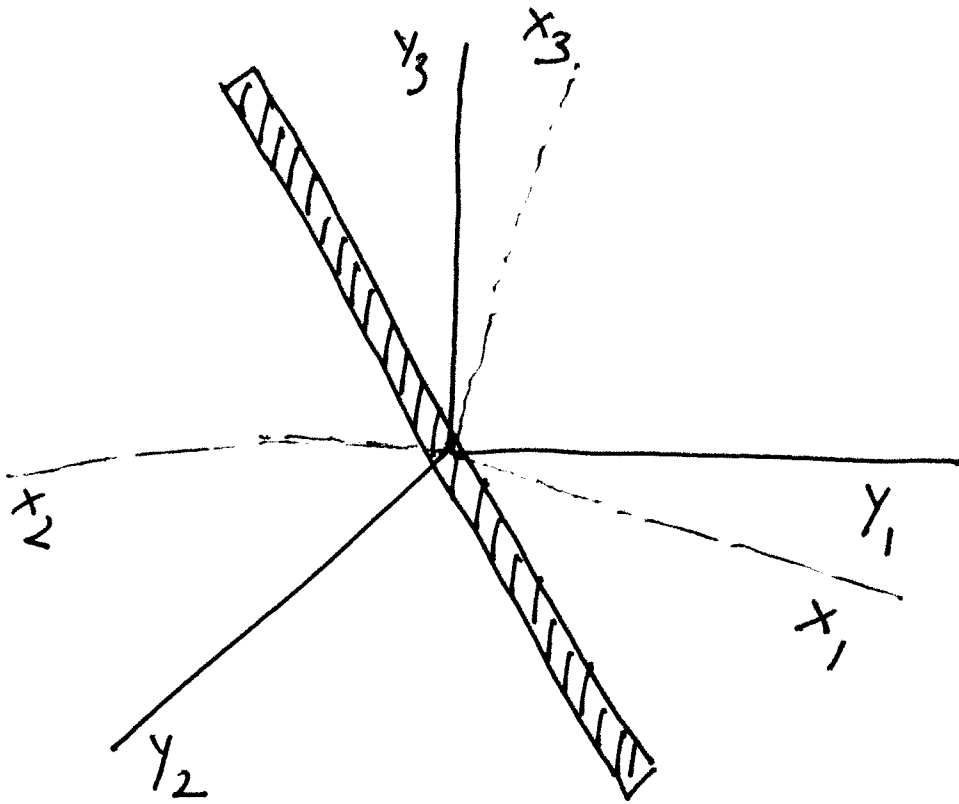
Hence

$$\begin{pmatrix} \dot{\bar{\omega}}_1 \\ \dot{\bar{\omega}}_2 \\ \dot{\bar{\omega}}_3 \end{pmatrix} = \begin{pmatrix} 0 & \alpha & 0 \\ -\alpha & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \bar{\omega}_1 \\ \bar{\omega}_2 \\ \bar{\omega}_3 \end{pmatrix}$$



④

Consider the problem of orienting one set of co-ordinate axes w.r.t. a second set



Say that the projection on the  $y_j$  axis of a unit vector along the  $x_i$  axis is  $a_{ij}$ . These  $a_{ij}$ -s are called direction cosines and we can put them in the form of a matrix  $A$ .

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There are nine such direction cosines.

If the  $x$ -system is rotating about  
 its  $x_1$  axis with an angular velocity  $\omega_1$   
 $x_2$  " " " " " "  $\omega_2$   
 $x_3$  " " " " " "  $\omega_3$

Then  $A(t)$  changes with time as follows:

$$\dot{A}(t) = \underline{\Omega}(t) A(t)$$

$$\underline{\Omega} = \begin{pmatrix} 0 & \omega_3 & -\omega_2 \\ -\omega_3 & 0 & \omega_1 \\ \omega_2 & -\omega_1 & 0 \end{pmatrix}$$

— x —

Thus the orientation of a symmetrical rigid body in torque free spinning motion is equivalent to finding the transition matrix for the system

(6)

$$\dot{x}(t) = \Omega(t)x(t) \quad (*)$$

where

$$\Omega(t) =$$

$$\begin{pmatrix} 0 & \omega_0 & \omega_0' \sin \alpha t \\ -\omega_0 & 0 & \omega_0' \cos \alpha t \\ -\omega_0' \sin \alpha t & -\omega_0' \cos \alpha t & 0 \end{pmatrix}$$

where  $\omega_0, \omega_0', \alpha$  are all constants and

$$\alpha = \frac{I_2 - I_3}{I_2} \omega_0$$

Problem: Find the state transition matrix of (\*).